

## 公募研究：ワニエ関数を軸とする準粒子自己無撞着法の新しい展開

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### 第一原理計算の基礎的な方法論

一体問題→低エネルギー励起、準粒子の寿命など  
→不純物や分子、表面などへの応用

- \* 一体問題解法を整備する(オーダーNではない範囲)。局在基底と非局在基底。
- \* 「準粒子self-consistent GW法」の整備  
(異方性、取り扱いやすさ、収束性、計算速度)。
- \* ワニエ関数の構成。低エネルギー励起(フォノンやマグノンなど)。  
それらによる自己エネルギーなど。

# The **PMT** method: a new linearized method

$$\text{PMT} = \text{APW} + \text{MTO}$$

Muffin-tin

*Takao Kotani (tottori-u, Japan)*

## Outline

### 1. Introduction

Motivation, linearized method, APW, MTO

### 2. Formalism A few key points

### 3. Results

from  $\text{H}_2$  through  $\text{Kr}_2$  in PBE

ecalj

google

Not organized well yet ...  
(MostofAll are still in Japanese...)

# Motivation:

Quasiparticle Self-consistent GW (QSGW) method

$$\hat{H} = \hat{H}_0 + (\hat{H} - \hat{H}_0) \text{ と分ける方法}$$

をFP-LMTOをもとに開発してきた。

→ 方法論的には成功した。現状では  
数値的に問題点は色々ある。

Empty spheres,

Not systematic convergence,

計算のパラメータ設定が面倒。

→ Look for a better linearized method.

# Look for a better linearized method!

## Requirements:

### \* 速度、精度、ロバストネス:

Reliable in any atomic configurations.  
Stable convergence (intrinsic stability).  
Systematic convergence check.

### \* 更なる発展:

GW, Wannier functions, magnon, phonons...  
excitations, their coupling, and total energy.

## My conclusion:

To overcome shortcomings in FP-LAPW and FP-LMTO, we should **use both of APW and MTO together.**

# The linearized method

- **Basis set**

→ eigenfunctions are given as:

$$\psi_p(\mathbf{r}) = \sum_j \alpha_p^j F_j(\mathbf{r}) \quad \{F_j\} : \text{basis set}$$

- **Total energy in DF as a functional of  $\{\psi_p\}$ .**



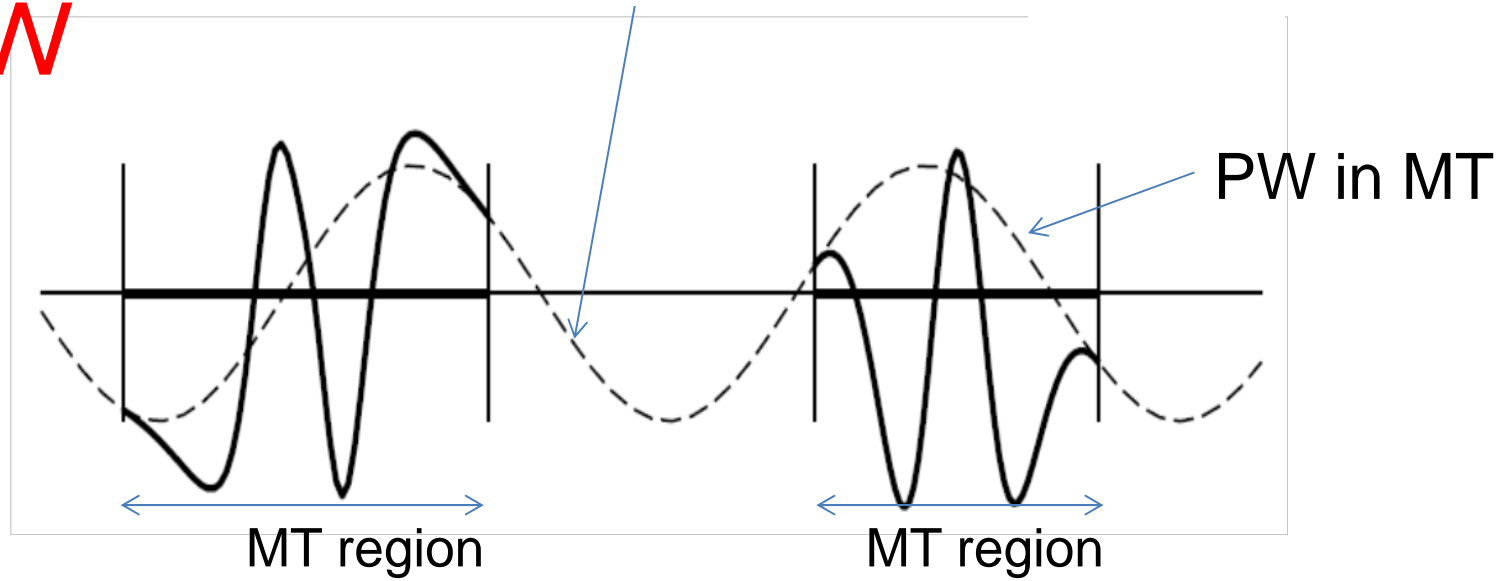
Kohn-Sham equation

- For fixed  $\{F_j\}$ , minimization with respect to  $\{\alpha_p^j\}$ .

- For force, we take into account the changes of  $\{F_j\}$ ; the Pulley term.

Envelope function = Plane Wave(PW)

APW

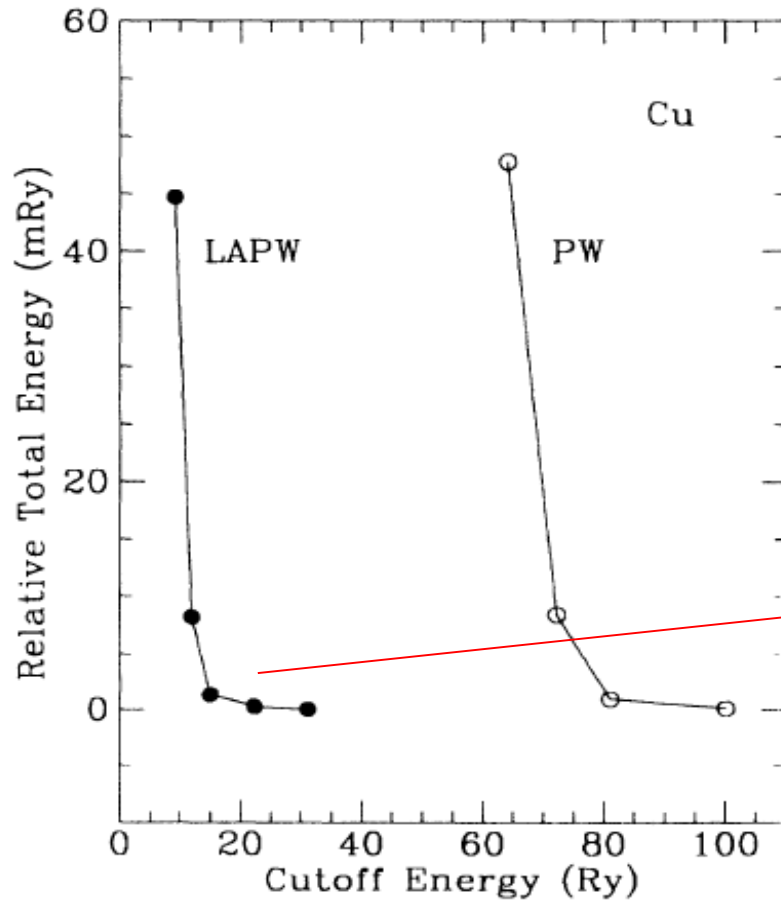


### 3-components

PW + Atomic-like parts — Counter parts  
0<sup>th</sup> 1<sup>st</sup> 2<sup>nd</sup>  
within MTs within MTs

- Good for extended state such as Na(3s).
- Systematic
- Inefficient for Cu(3d), O(2p)

# Linearized APW



fcc Cu

plane wave is not so efficient for 3d

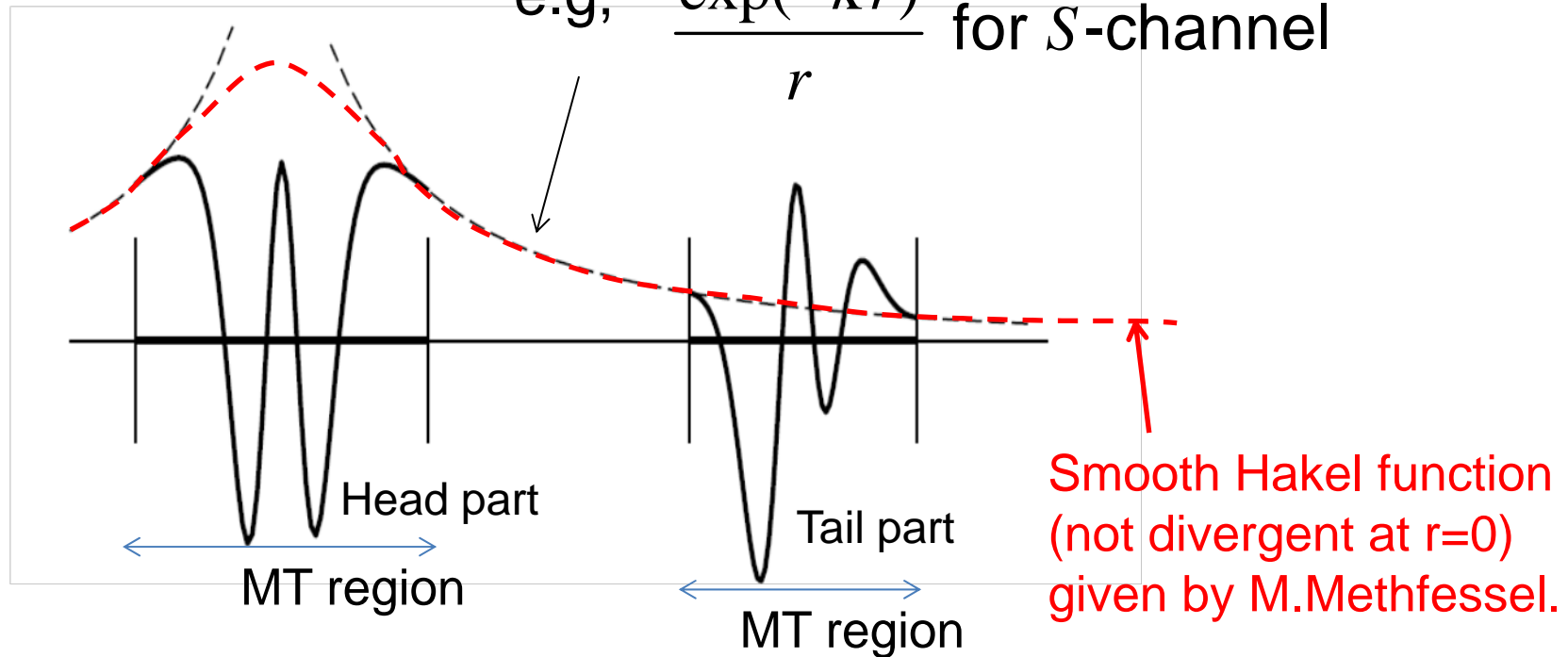
15Ry or more for ~1mRy convergence

FIG. 1. Relative total energy of fcc Cu plotted against plane-wave cutoff energy.

D.Singh et al PRB49,17424

# MTO

“Atom-centered Hankel function”  $\times Y_{lm}$ ,  
e.g,  $\frac{\exp(-\kappa r)}{r}$  for  $S$ -channel



## 3-components

Hankel + Atomic-like parts — Counter parts of Hankel  
within MT within MT

0<sup>th</sup>

1<sup>st</sup>

2<sup>nd</sup>

- Good for localized orbitals such as Cu(3d), O(2p)
- Not so good for extended states.
- Not systematic



$$\text{PMT} = \text{APW} + \text{MTO}$$

This is physically reasonable. Complementary.

*T.K and M.van Schilfgaarde*

arXiv:0808.1604

Phys. Rev. B 81, 125117 (2010)

→ Highly accurate for solids.

No empty sphere.

Small number of APWs.

(Parameters of MTO controlled by hand).

--- Recent development ---

*T.K and Hiori Kino*

- Formulation based on a general theory of augmentation.

- Results for dimers from H<sub>2</sub> through Kr<sub>2</sub> in PBE.

Simple setting of MTOs with APWs (Energy cutoff 2~4Ry) .

# 2. Formalism

A general formalism of augmentation wave.

Key points:

- 3-component Hilbert space
- Augmentation for product of envelope functions.

# A basis function

- MT centers at  $\mathbf{R}_a$ . Radius  $R_a$ .
- Spin index is not shown explicitly.

a basis  $F_j(\mathbf{r})$  consists of 3-kinds of components

$$F_{0j}(\mathbf{r}), \quad \left\{ F_{1j,a}(\mathbf{r}) \right\}, \quad \left\{ F_{2j,a}(\mathbf{r}) \right\}$$

Envelope function

Atomic-like parts

Counter parts

PW or

for  $|\mathbf{r}| < R_a$

for  $|\mathbf{r}| < R_a$

Hankel(Bloch sum)

Cutoff:  $l \leq l_{\max} \sim 4$ , Radial-part :  $\sim 5^{\text{th}}$ -order Laguerre polynomials

Augmentation procedure depend on Envelop functions

Conventional representation of augmented wave (Soler-Williams type)

$$F_j(\mathbf{r}) = F_{0j}(\mathbf{r}) + \sum_a F_{1j,a}(\mathbf{r} - \mathbf{R}_a) - \sum_a F_{2j,a}(\mathbf{r} - \mathbf{R}_a)$$

## Conventional Soler-Williams formalism

$$F_j(\mathbf{r}) = F_{0j}(\mathbf{r}) + \sum_a F_{1j,a}(\mathbf{r} - \mathbf{R}_a) - \sum_a F_{2j,a}(\mathbf{r} - \mathbf{R}_a)$$

● Eigenfunctions  $\psi_p = \sum_j \alpha_p^j F_j$

● Product 
$$F_i^*(\mathbf{r})F_j(\mathbf{r}) = \left( F_{0i}^*(\mathbf{r}) + \sum_a F_{1i,a}^*(\mathbf{r} - \mathbf{R}_a) - \sum_a F_{2i,a}^*(\mathbf{r} - \mathbf{R}_a) \right) \times \left( F_{0j}(\mathbf{r}) + \sum_a F_{1j,a}(\mathbf{r} - \mathbf{R}_a) - \sum_a F_{2j,a}(\mathbf{r} - \mathbf{R}_a) \right)$$

contains cross terms. → Approximaiton: we neglect these terms  
→ DF total energy, Kohn-Sham equation.

A problematic logic...

It is better to say **no cross terms by definition.**

How to formulate?(next page)

we start over from the definition of the Hilbert space.

## 3-component Hilbert space

A basis itself is a set of functions :

$$F_j = \left\{ F_{0j}(\mathbf{r}), \left\{ F_{1j,a}(\mathbf{r}) \right\}, \left\{ F_{2j,a}(\mathbf{r}) \right\} \right\} \quad \{\dots\} \text{ means a set.}$$

The set of  $F_j$ ,  $\{F_j\}$  is a Hilbert space  
(as a direct sum of linear spaces).

→ 3-component Hilbert space

We define a quantum mechanical system in this space. Then we apply DFT.

In the 3-component Hilbert space

$$F_j = \left\{ F_{0j}(\mathbf{r}), \left\{ F_{1j,a}(\mathbf{r}) \right\}, \left\{ F_{2j,a}(\mathbf{r}) \right\} \right\},$$

• Eigenfunction:  $\psi_p = \sum_j \alpha_p^j F_j$

• Definition of product for bilinear quantities:

$$F_i^*(\mathbf{r})F_j(\mathbf{r}') = \left\{ F_{0i}^*(\mathbf{r})F_{0j}(\mathbf{r}'), \left\{ F_{1i,a}^*(\mathbf{r})F_{1j,a}(\mathbf{r}') \right\}, \left\{ F_{2i,a}^*(\mathbf{r})F_{2j,a}(\mathbf{r}') \right\} \right\}$$

$$\psi_{p'}^*(\mathbf{r})\psi_p(\mathbf{r}') = \sum_{i,j} \alpha_{p'}^{i*} \alpha_p^j F_i^*(\mathbf{r})F_j(\mathbf{r}')$$

Product is defined  $\rightarrow$  How to connect to real world?

**Mapping from the 3-component Hilbert space to real world:**

$$F_i^*(\mathbf{r})F_j(\mathbf{r}') = \left\{ F_{0i}^*(\mathbf{r})F_{0j}(\mathbf{r}'), \left\{ F_{1i,a}^*(\mathbf{r})F_{1j,a}(\mathbf{r}') \right\}, \left\{ F_{2i,a}^*(\mathbf{r})F_{2j,a}(\mathbf{r}') \right\} \right\}$$

**$\mathcal{R}$ -mapping,  $\mathcal{R}[\dots]$**

$$\mathcal{R}\left[ F_i^*(\mathbf{r})F_j(\mathbf{r}') \right] = F_{0i}^*(\mathbf{r})F_{0j}(\mathbf{r}') + \sum_a F_{1i,a}^*(\mathbf{r}-\mathbf{R}_a)F_{1j,a}(\mathbf{r}'-\mathbf{R}_a) - \sum_a F_{2i,a}^*(\mathbf{r}-\mathbf{R}_a)F_{2j,a}(\mathbf{r}'-\mathbf{R}_a)$$

Meaningful when  $\mathbf{r} \approx \mathbf{r}'$ :

(1) Both are in the same MT; (2) Both are in interstitial region.

**augmentation for the product of Envelope functions**  
 (the essence of Soler-Williams formulation)

c.f.

$$\mathcal{R}\left[ F_i^*(\mathbf{r}) \right] \mathcal{R}\left[ F_j(\mathbf{r}') \right]$$

$\mathcal{R}$ -mapping is not to basis functions, but to the products.

“density-matrix-like product”

$$\psi_p = \sum_j \alpha_p^j F_j \quad \rho_{ij} = \sum_P^{\text{OCCUPIED}} \alpha_p^{i*} \alpha_p^j$$

$$\sum_{i,j} \rho_{ij} F_i^*(\mathbf{r}) F_j(\mathbf{r}') = \left\{ \begin{array}{l} \sum_{i,j} \rho_{ij} F_{0i}^*(\mathbf{r}) F_{0j}(\mathbf{r}'), \\ \sum_{i,j} \rho_{ij} F_{1i}^*(\mathbf{r}) F_{1j}(\mathbf{r}'), \quad \sum_{i,j} \rho_{ij} F_{2i}^*(\mathbf{r}) F_{2j}(\mathbf{r}') \end{array} \right\} \quad \text{3-component}$$

↓  $\mathcal{R}$ -mapping

$$\mathcal{R} \left[ \sum_{i,j} \rho_{ij} F_i^*(\mathbf{r}) F_j(\mathbf{r}') \right] = \sum_{i,j} \rho_{ij} F_{0i}^*(\mathbf{r}) F_{0j}(\mathbf{r}')$$

$$+ \sum_{i,j} \rho_{ij} F_{1i}^*(\mathbf{r} - \mathbf{R}) F_{1j}(\mathbf{r}' - \mathbf{R}) - \sum_{i,j} \rho_{ij} F_{2i}^*(\mathbf{r} - \mathbf{R}) F_{2j}(\mathbf{r}' - \mathbf{R})$$

→ density, kinetic energy density, and so on in the Hilbert space.

(Another approximation to the Hartree term (Weinert method))

→ as a functional of the 3-component density  $n(\mathbf{r}) = \sum_{i,j} \rho_{ij} F_i^*(\mathbf{r}) F_j(\mathbf{r})$



Finally, we can give total energy in DFT as a functional of  $\psi_p = \sum_j \alpha_p^j F_j$



Minimization of the total energy gives the Kohn-Sham equation.

Our formalism results in the equivalent equations with previous articles, but :

- Consistency is transparent.
- Definition of diatomic quantities in the cutoff model is clear.
- Convenient to derive atomic forces (clear about surface term).

## Hilbert space in the Real-world

A basis function

$$F_i(\mathbf{r}) = F_{0i}(\mathbf{r}) + F_{1i}(\mathbf{r} - \mathbf{R}) - F_{2i}(\mathbf{r} - \mathbf{R})$$

product  $F_i^*(\mathbf{r})F_j(\mathbf{r}')$

**Bilinear quantities:**  
norm, density, and so on.  
**Coulomb interaction also.**

No cutoff limit

Exact Solution

## 3-component Hilbert space (artificial)

A basis function

$$F_i(\mathbf{r}) = \{F_{0i}(\mathbf{r}), F_{1i}(\mathbf{r}), F_{2i}(\mathbf{r})\}$$

**Define** product  $F_i^*(\mathbf{r})F_j(\mathbf{r}')$

By definition, No cross terms

**Define bilinear quantities,**  
**Coulomb interaction and**  
**so on.**

No cutoff limit

Exact Solution

Both should be the same

Basis is in real world → Approximation

Numerically solvable model.  
Well-defined for finite cutoff.

*c.f.* PAW treated the Hilbert space of “pseudo wave functions”.  
(  $F_{0i}(\mathbf{r})$  determines  $F_{1i}(\mathbf{r})$  and  $F_{2i}(\mathbf{r})$  uniquely. )

- Prepare 3-components Hilbert space
  - Define products and so on in the space
  - For given  $\psi_p = \sum_j \alpha_p^j F_j$ ,  
we can evaluate<sup>j</sup> (or define) bilinear quantities  
as a functional in the space through  **$\mathcal{R}$ -mapping**.
- 

Problems:

Positive definiteness of overlap matrix

Null vector problems

Orthogonality to core (since we use frozen core approximation).

# RESULTS

- PBE, Colinear Spin, Scalar Relativistic DFT
- Large supercell  $13.5 \text{ \AA} \times 15 \text{ \AA} \times 16.5 \text{ \AA}$ .
- APW and Atom-centered localized basis (MTO and lo)
- Ferro-magnetic. Only inversion symmetry.  $\Gamma$  point only.

MTO set (part)	Mg	2s2p(lo)	$3s3p3d4f$	$\times 2$
	Al	2s2p(lo)	$3s3p3d4f$	
	Si		$3s3p3d4f$	
	P		$3s3p3d4f$	
	S		$3s3p3d4f$	
	Cl		$3s3p3d4f$	
	Ar	3s3p(lo)	$4s4p3d4f$	
	K	3s3p(lo)	$4s4p3d4f$	
	Ca	3s3p(lo)	$4s4p3d4f$	
	Sc	3s3p(lo)	$4s4p3d4f$	
	Ti	3s3p(lo)	$4s4p3d4f$	

Number of local basis  $(1+3+5+7) \times 2 + (1+3) = 36$

Atom-independent setting of MTO parameters

(not shown here)

# Atom in supercell

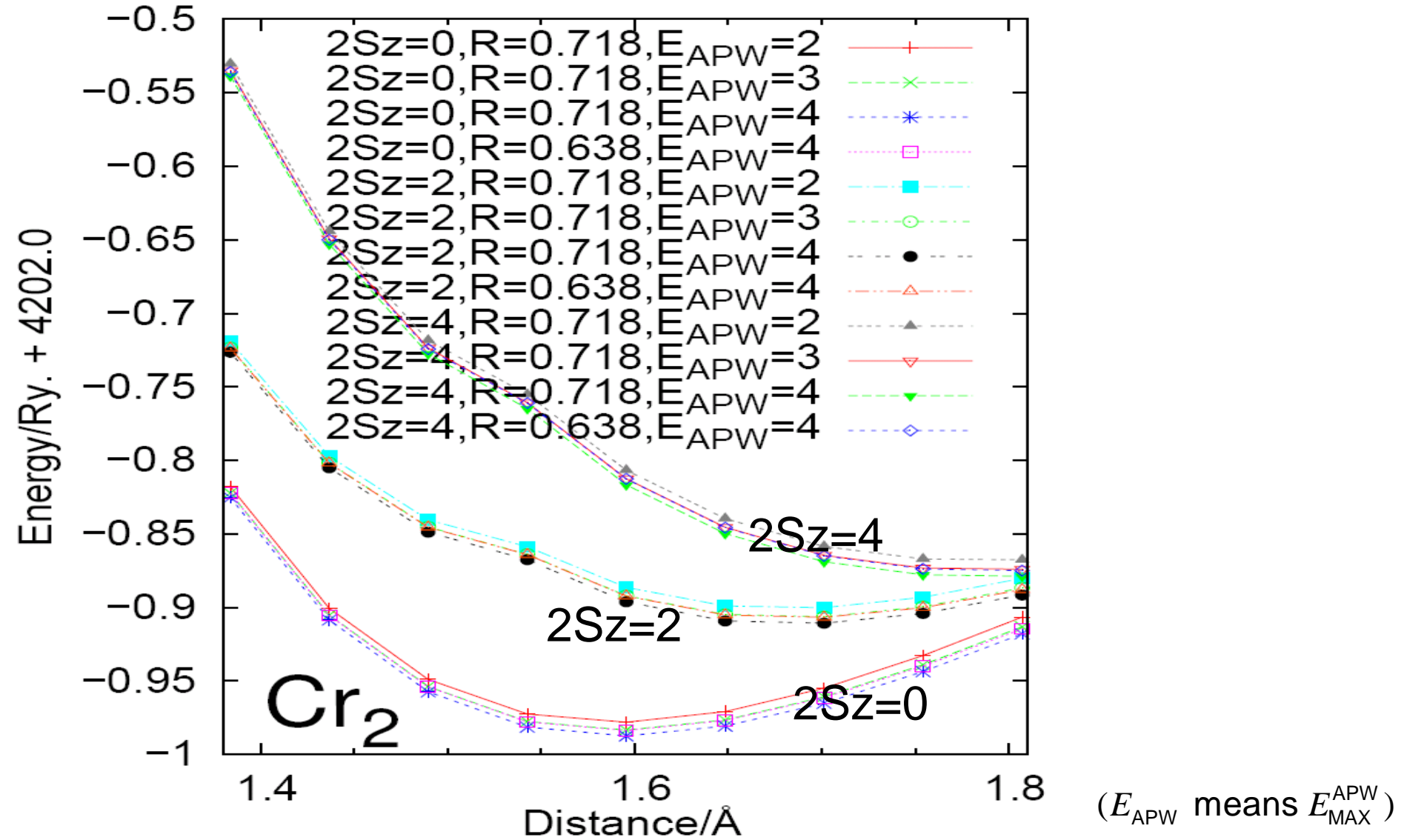
	$E_{\text{MAX}}^{\text{APW}} = 2$	$E_{\text{MAX}}^{\text{APW}} = 3$	$E_{\text{MAX}}^{\text{APW}} = 4$	(Ry)
$0, 2S_z = 2$	$E_e/\text{Ry} = -150.0 + \Delta E_e/\text{Ry}$			
$R = 0.427 \text{ \AA}$	-0.10930,	-0.12201,	-0.12852	
$R = 0.488 \text{ \AA}$	-0.12950,	-0.13640,	-0.13949	
$R = 0.549 \text{ \AA}$	-0.13609,	-0.13983,	-0.14139	
$R = 0.610 \text{ \AA}$	-0.13913,	-0.14098,	-0.14172	
$\text{Cr}, 2S_z = 6$	$E_e/\text{Ry} = -2101.0 + \Delta E_e/\text{Ry}$			
$R = 0.638 \text{ \AA}$	-0.42852,	-0.43596,	-0.44023	
$R = 0.718 \text{ \AA}$	-0.43316,	-0.43807,	-0.44141	
$R = 0.797 \text{ \AA}$	-0.43110,	-0.43763,	-0.44198	

$$N_{\text{APW}} = 1081, \quad 1973, \quad 3025$$

$$N_{\text{BASIS}} = 1081 + 36, \quad 1973 + 36, \quad 3025 + 36$$

# Diatomic molecules

(From H<sub>2</sub> through Kr<sub>2</sub>)



	$E_{APW}=2$	$E_{APW}=3$	$E_{APW}=4$
$N_{APW}$	1081	1973	3025
$N_{BASIS}$	1081+36x2	1973+36x2	3025+36x2

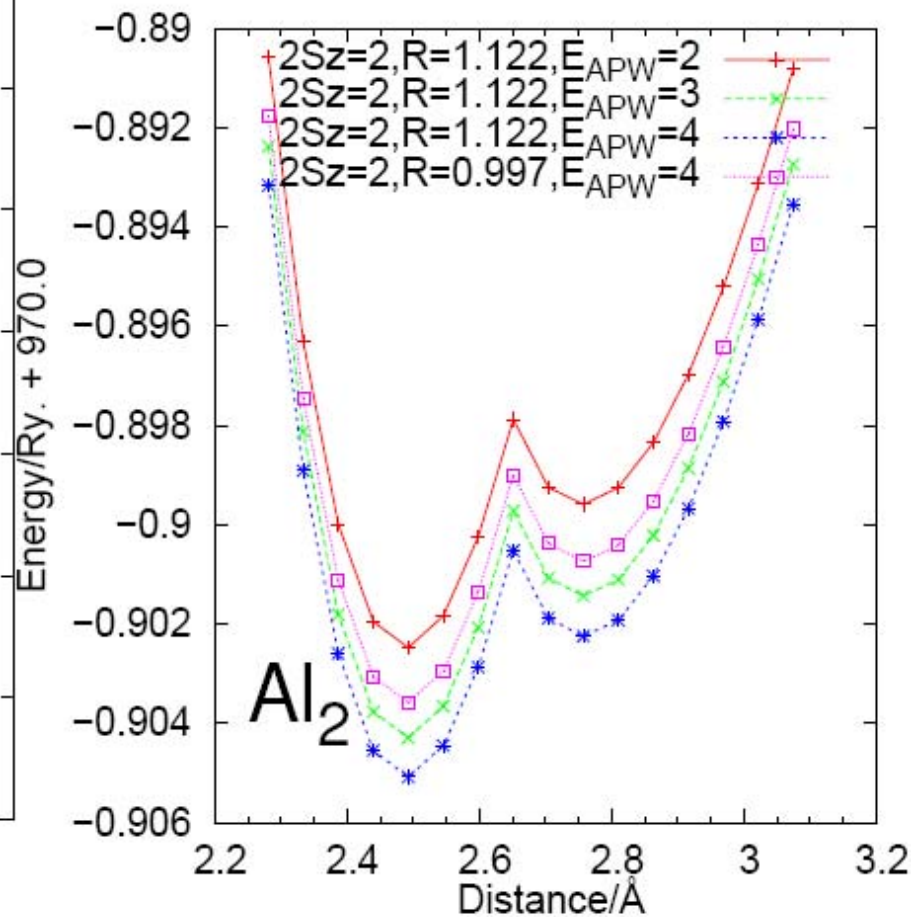
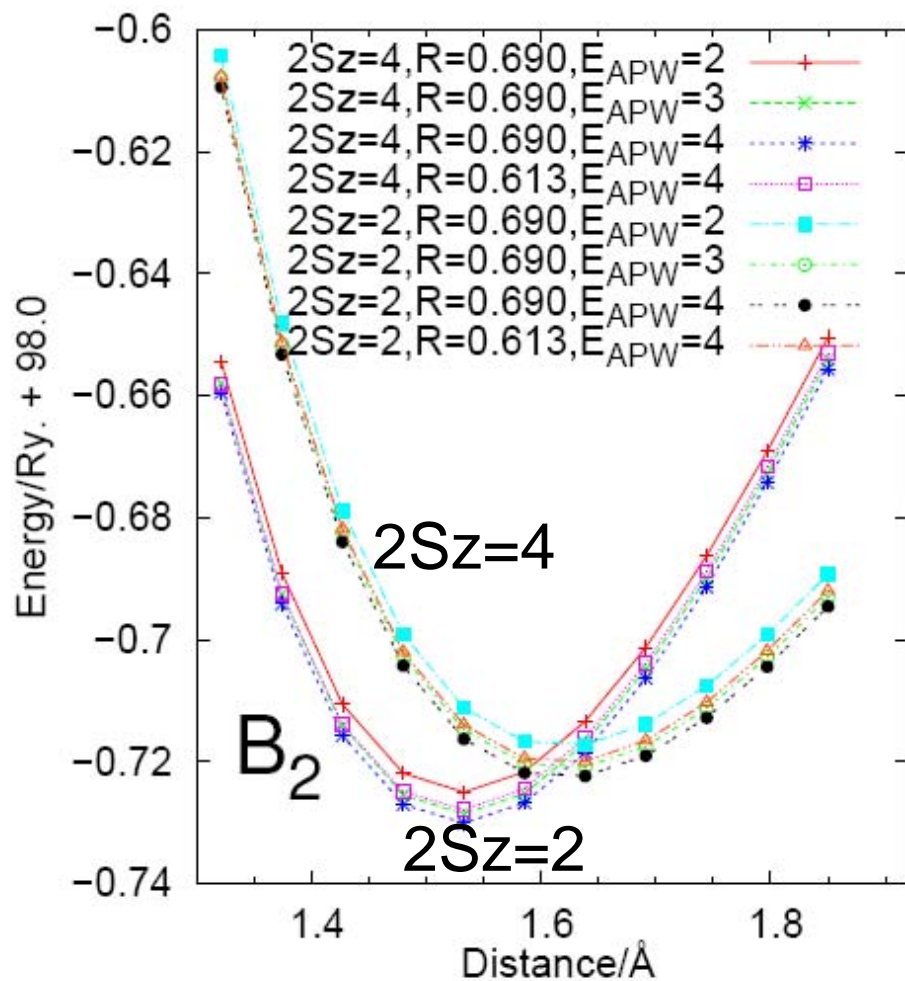
# Cr<sub>2</sub> at equilibrium atomic distance

$E_{\text{MAX}}^{\text{APW}} / \text{Ry}$ ( $N_{\text{APW}}$ )	Equilibrium Atomic distance			Vibrational frequency
	$r_e / \text{\AA}$	$E_e / \text{Ry}$	$D_e / (\text{Kcal/mol})$	$\omega_e / (\text{cm}^{-1})$
2 (1081)	1.589	-4202.97816	35.0868	819.273
3 (1973)	1.590	-4202.98356	33.7052	816.532
4 (3025)	1.591	-4202.98747	32.8334	813.296
5 (4245)	1.591	-4202.99086	32.2915	811.480
6 (5573)	1.592	-4202.99388	31.9939	808.289
7 (7025)	1.593	-4202.99664	31.8474	806.676
8 (8611)	1.593	-4202.99917	31.8164	806.530

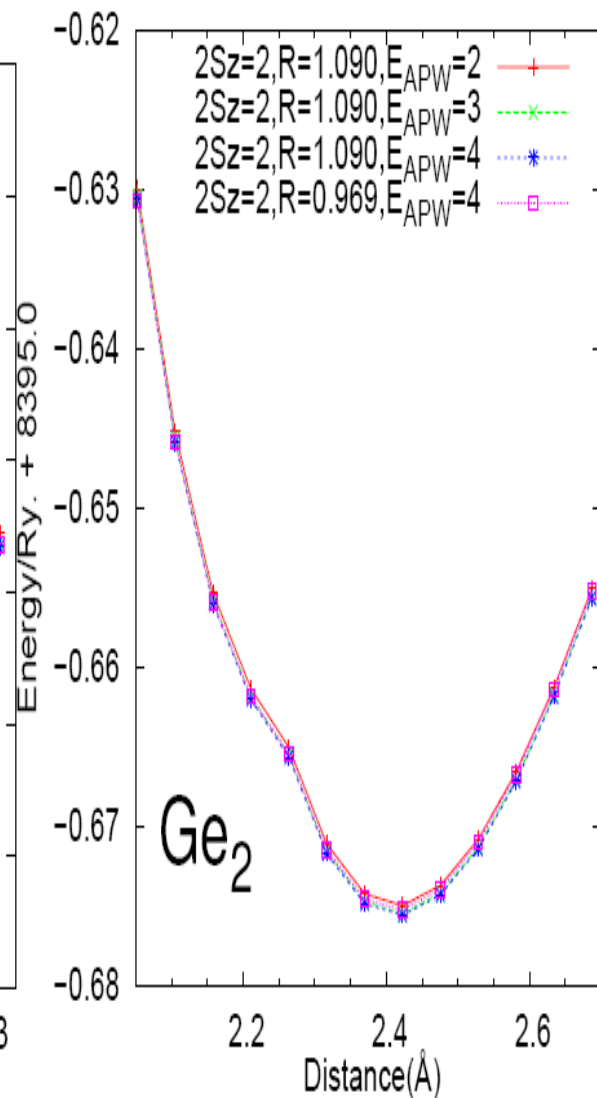
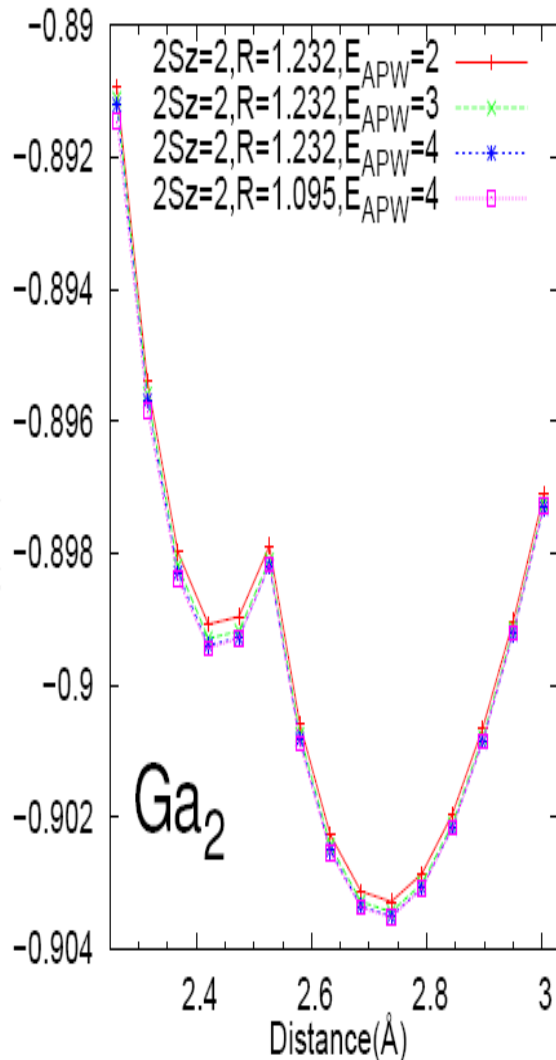
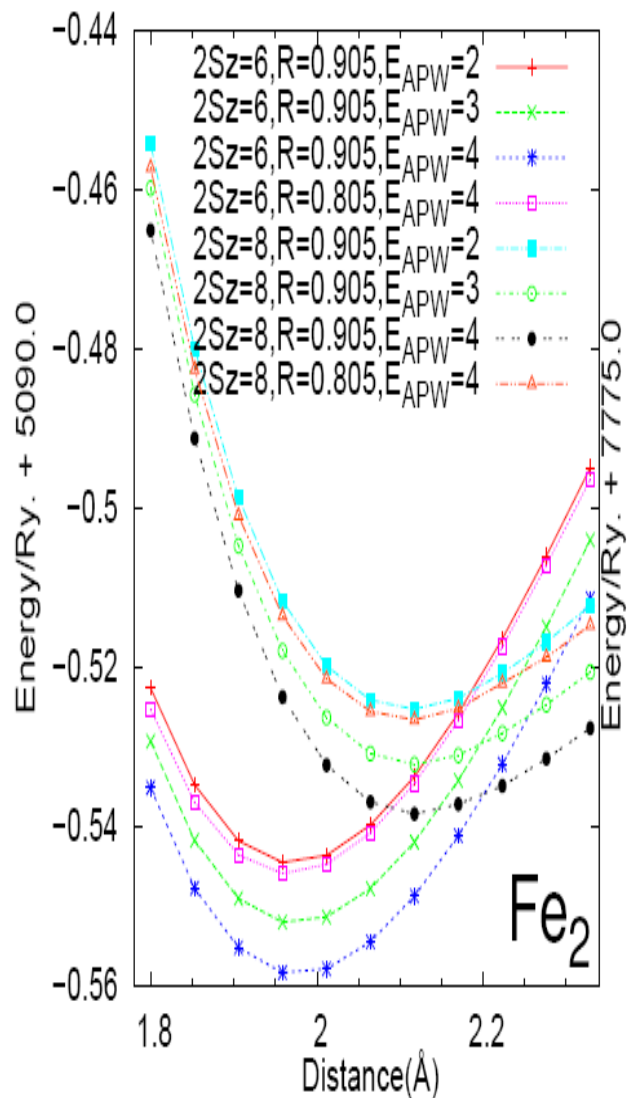
Atomization energy

$1\text{eV} = 1/13.605 \text{ Ry} = 23.06 \text{ Kcal/mol.}$





$$N_{\text{basis}} = 1081 + 36 \times 2, \quad 1973 + 36 \times 2, \quad 3025 + 36 \times 2$$



$$N_{\text{basis}} = 1081 + 36 \times 2, \quad 1973 + 36 \times 2, \quad 3025 + 36 \times 2$$

## Comparison with Gaussian

		$r_e$ (Å)	$D_e$ (Kcal/mol)	$\omega_e$ (cm <sup>-1</sup> )
$H_2, 2S_z=0$	PMT	0.749	104.678	4317.959
	PMT(NR)	0.750	104.764	4311.202
	GTO	0.752	104.552	4311.816
$O_2, 2S_z=2$	PMT	1.218	143.741	1564.787
	PMT(NR)	1.218	144.984	1568.867
	GTO	1.220	<del>139.815</del>	1554.249
	VASP		143.3	
$Cr_2, 2S_z=0$	PMT	1.591	32.833	813.296
	PMT(NR)	1.589	30.191	818.483
	GTO	1.595	26.192	808.148
$Fe_2, 2S_z=6$	PMT	1.977	57.596	397.673
	PMT(NR)	1.991	58.770	386.597
	GTO	2.012	56.902	397.228
$Cu_2, 2S_z=0$	PMT	2.218	51.169	269.326
	PMT(NR)	2.251	48.503	254.321
	GTO	2.251	48.645	255.768

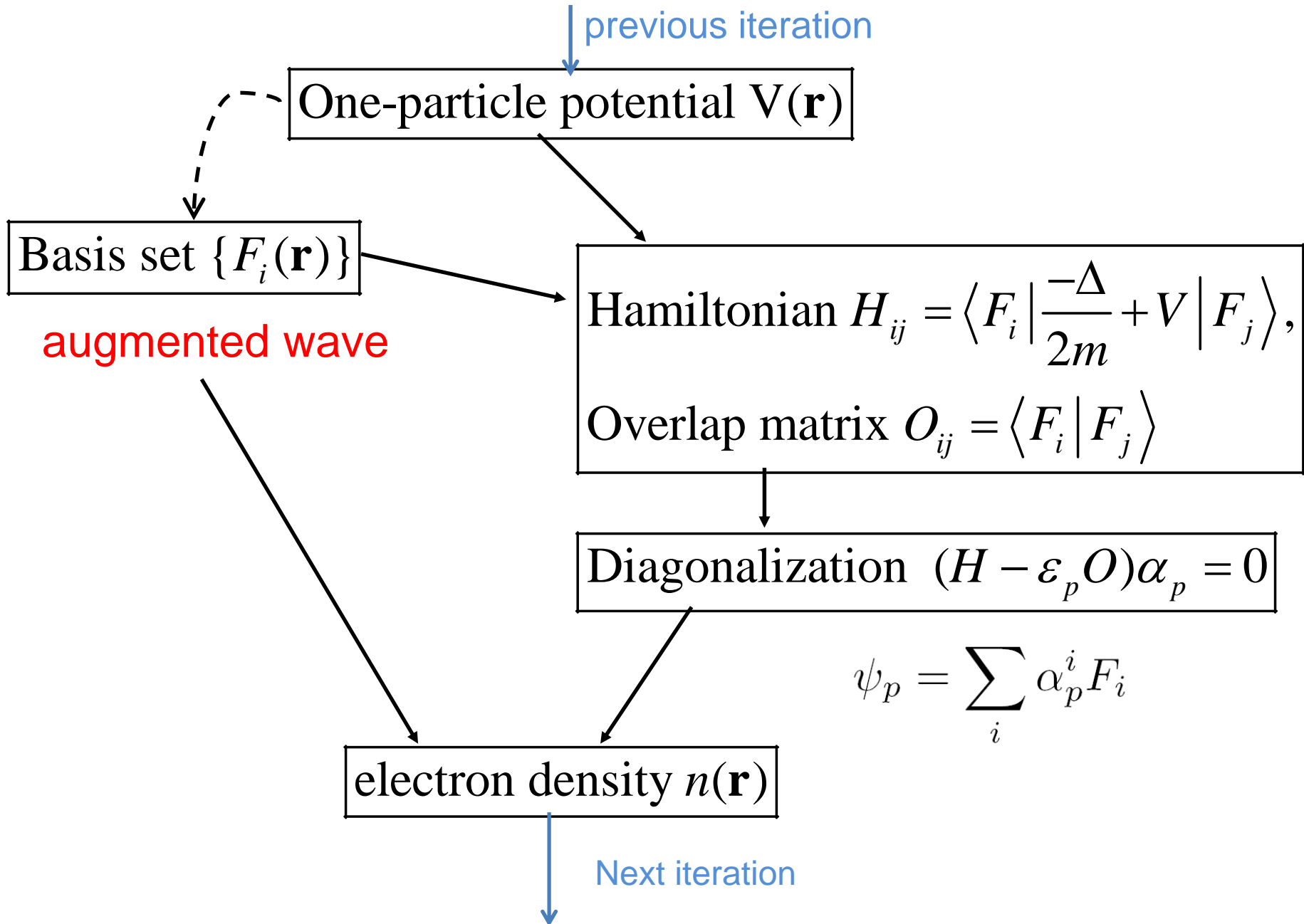
NR: non relativistic      GTO: 6-311+G(d,p)

# Summary

- Formalism based on “3-component Hilbert space”.  
Theoretical consistency is transparent for cutoffted model.  
Atomic force is derived systematically (not shown).
- With APW (cutoff energy is  $2 \sim 4\text{Ry}$ ), we can perform accurate calculations for homo nuclear diatomic molecules.
- Two kinds of augmented wave  $\rightarrow$  very efficient calculations.  
For example, 300 atoms in  $15 \text{ \AA}$  cell  $\rightarrow$  roughly speaking,  
basis  $\sim 30 \times 300 + 2000$  ( $\sim 10000$  basis)
- \* Construction of the Wannier functions. (superposition of MTO).
- \* Better control for Positive definiteness of Norm.  
Null space problem:  
vectors with zero eigenvalues for Hamiltonian and Overlap matrix.



# Total energy minimization in the Linearized method



• inner product:  $\langle F_i | F_j \rangle =$

$$\int_{\text{Cell}} d^3 r F_{0i}^*(\mathbf{r}) F_{0j}(\mathbf{r}) + \int_{|\mathbf{r}| < R} d^3 r F_{1i}^*(\mathbf{r}) F_{1j}(\mathbf{r}) - \int_{|\mathbf{r}| < R} d^3 r F_{2i}^*(\mathbf{r}) F_{2j}(\mathbf{r})$$

$$\int_{\text{Cell}} d^3 r F_{0i}^*(\mathbf{r}) F_{0j}(\mathbf{r}) - \int_{|\mathbf{r}| < R} d^3 r F_{2i}^*(\mathbf{r}) F_{2j}(\mathbf{r})$$

should be positive definite. Also for the Coulomb matrix.

## Hilbert space in the Real-world

A basis function

$$F_i(\mathbf{r}) = F_{0i}(\mathbf{r}) + F_{1i}(\mathbf{r}-\mathbf{R}) - F_{2i}(\mathbf{r}-\mathbf{R})$$

Basis for product  $F_i^*(\mathbf{r})F_j(\mathbf{r}')$

**Bilinear Operators**

No cutoff limit

Exact Solution

## 3-component Hilbert space(model space)

A basis function

$$F_i(\mathbf{r}) = \{F_{0i}(\mathbf{r}), F_{1i}(\mathbf{r}), F_{2i}(\mathbf{r})\}$$

Basis for product  $F_i^*(\mathbf{r})F_j(\mathbf{r}')$  No cross terms

**Define bilinear operators in this space**

No cutoff limit

Exact Solution

$\mathcal{R}$ -mapping

**Both should be the same**

Approximations to  $\rightarrow$  Solve it

Numerically solvable model

1. Numerically solvable model  $\rightarrow$  then apply DFT
2. The exact solution at the infinite cutoff limit.

*c.f.* PAW gives “quantum mechanism” in space of “pseudo wave functions”.



## Hilbert space in the Real-world

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